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## Computational Mechanics

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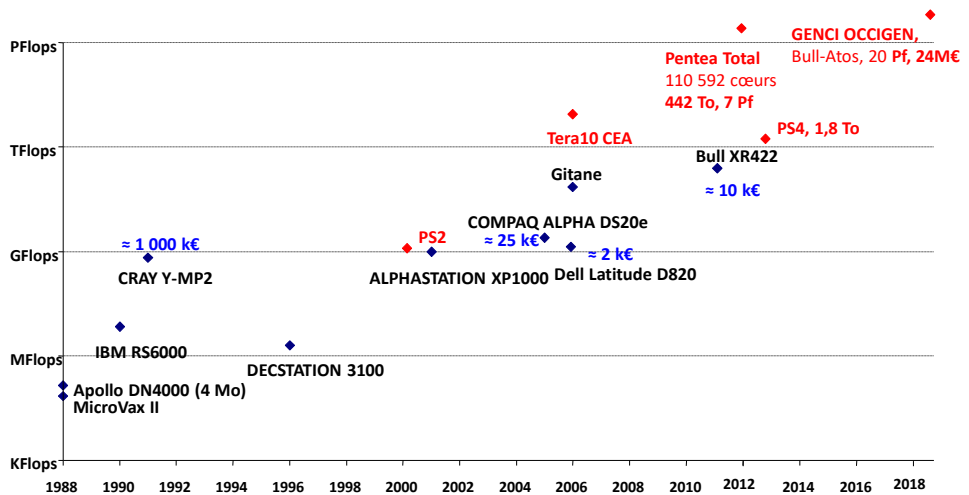





- Computational Mechanics
- Numerical simulation process
- Finite Difference Method
- Variational approach
- Finite Element Method
- Finite Volume Method
- Galerkin Discontinuous Finite Element Method

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### Computational Mechanics

Evolution of computeurs





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### Computational Mechanics

#### Evolution of computational mechanics

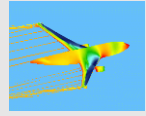
Aim:

- Reproduce to explain the inaccessible
- Plan before you realize

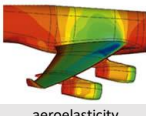
Replace  
Theory-Experiment  
by  
Theory-computation-Experiment

Evolution:

- Multiphysics
- Multiscale
- High Performance Computing



aerodynamic



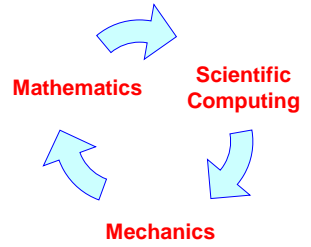
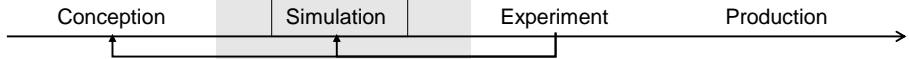

aeroelasticity



Structural analysis



Sources : engineering.swan  
ONERA & Lecture NF04 UTC

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### Computational Mechanics

#### Continuum mechanics

##### Conservation of mass

$$\frac{dp}{dt} + \rho \operatorname{div} \vec{v} = 0$$

or 
$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \vec{v}) = 0$$

##### Balance of linear momentum

$$\begin{cases} \operatorname{div} \vec{\sigma} + \vec{f} = \rho \vec{\gamma} & \text{dans } \Omega \\ \vec{\sigma} \vec{n} = \vec{F} & \text{sur } \partial\Omega \end{cases}$$

##### Balance of angular momentum

$$\vec{\sigma} = \vec{\sigma}^T$$

##### Conservation of energy

$$\rho \frac{de}{dt} - \vec{\sigma} : \vec{D} = r - \operatorname{div} \vec{q}$$

[m]  $\vec{x}$  : actual position

[s] t : time

$(\vec{x}, t)$ : Eulerian variables

[kg/m<sup>3</sup>]  $\rho(\vec{x}, t)$  : density

[m/s]  $\vec{v}(\vec{x}, t)$  : velocity

[N/m<sup>3</sup>]  $\vec{f}$  : body force per unit volume acting in  $\Omega$

[N/m<sup>2</sup>]  $\vec{F}$  : contact force per unit surface acting on  $\partial\Omega$

[−]  $\vec{n}$  : outward normal vector of  $\partial\Omega$

[m/s<sup>2</sup>]  $\vec{\gamma}$  : acceleration

[N/m<sup>2</sup>]  $\vec{\sigma}(\vec{x}, t)$  : Second order Cauchy stress tensor

[W/kg] e : specific internal energy

[s<sup>-1</sup>]  $\vec{D}$  : strain-rate (or stretching) tensor

[W/m<sup>2</sup>]  $\vec{q}$  : heat flux


[W/m<sup>3</sup>] r : heat supply density

Material derivative

$$\frac{d}{dt} = \frac{\partial}{\partial t} + (\vec{v} \cdot \nabla)$$

$$\vec{\gamma}(\vec{x}, t) = \frac{d\vec{v}}{dt} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v}$$

$$\vec{D} = \frac{1}{2} (\nabla^T \vec{v} + \nabla \vec{v})$$




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## Computational Mechanics

Example: Wave breaking

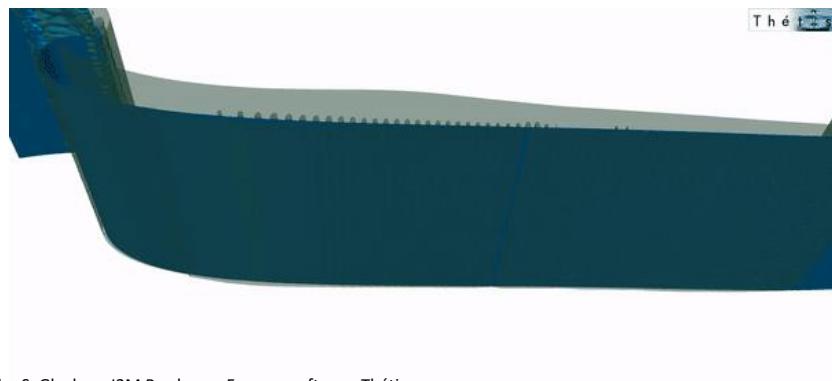
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## Computational Mechanics


Example: Wave breaking, full Navier-Stokes Air-Water simulation



P. Lubin, S. Glockner I2M Bordeaux, France, software Thétis

80 Millions cells on 1024 cores

800 Millions cells on 4096 cores



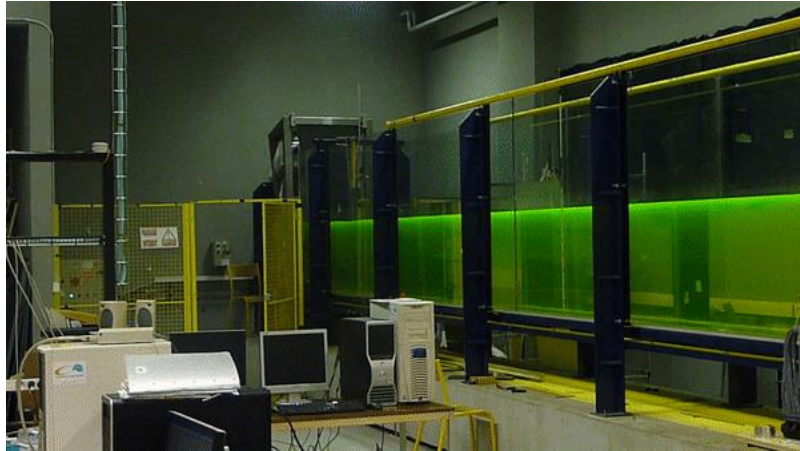
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## Computational Mechanics

Example: Wave propagation, wave breaking ..... experiment



Olivier Kimmoun, ECM, Marseille France, workshop B'Waves 2014 Bordeaux



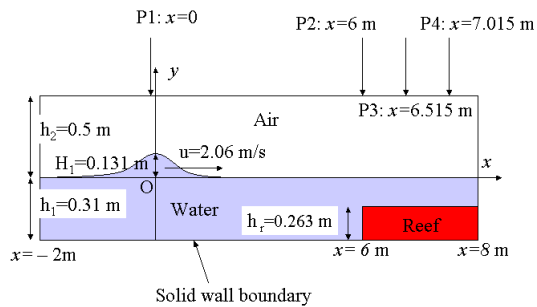
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## Computational Mechanics

Example : Wave propagation, wave breaking ..... experiment






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Computational Mechanics

Example: Wave propagation, wave breaking ..... Air-Water Euler model (Finite Volume)

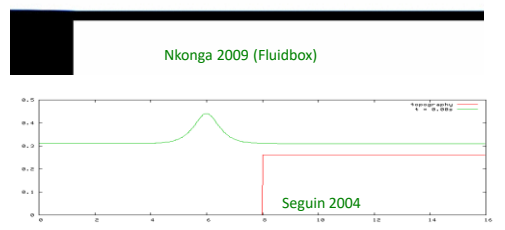
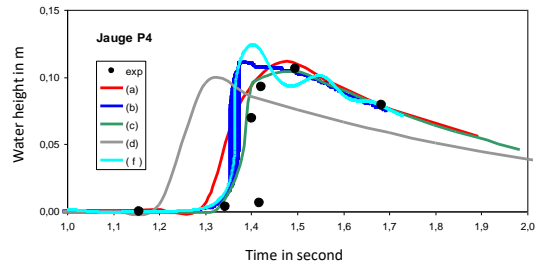
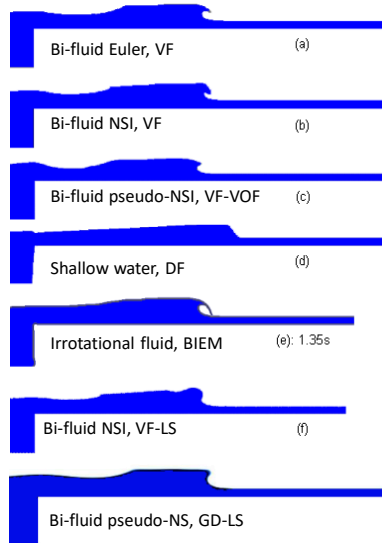



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Computational Mechanics

Confrontation of several CFD codes



P. Helluy, F. Golay, J.P. Caltagrone, P. Lubin, S. Vincent, D. Drevard, R. Marcer, P. Fraunie, N. Seguin, S. Grilli, A.N. Lesage, A. Dervieux, O. Allain, "Numerical simulation of wave breaking", M2AN, Vol.39 n°3, pp 591-607, 2005

B. Braconnier, J-J Hu, Y-Y Niu, B. Nkonga, K-M Shyue, "Numerical simulations of low Mach compressible two-phase flows: Preliminary assessment of some basic solution techniques", ESAIM: Proc., Vol. 28, pp. 117-134, 2009.




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## Computational Mechanics

Example, metal forming using finite elements (Forge3D – Cemef)


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## Computational Mechanics

Example of computation on GPU (DEM+CuDa)





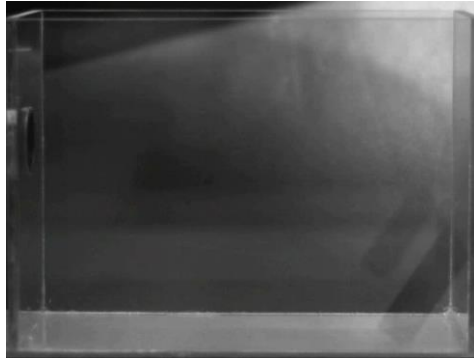
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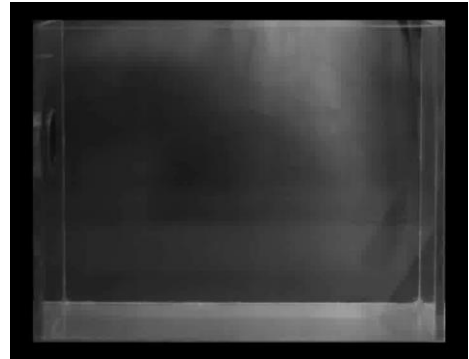
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## Computational Mechanics

### Realistic image rendering example applied to fluid mechanics




Experiment



Numerical simulation:  
FEM/VMS + LES +  
Convected Level Set...  
+ blender

Elie Hachem, Thierry Coupez, Cemef, code CimLib, anisotropic adaptive mesh refinement technic



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## Computational Mechanics

### Overview of numerical approximation of partial differential equations

- Finite Difference Method (FDM)
- Finite Element Method (FEM)
- Finite Volume Method (FVM)
- Spectral Method
- Boundary Element Method (BEM)
- Lattice Boltzmann (LB)
- Smooth Particle Hydrodynamics (SPH)
- Discrete Element Method (DEM)
- Many others ...

#### Finite Difference Method (FDM)

The oldest method, used for many physics

- Easy to implement
- Easy to make higher order
- Only applicable on structured grids
- Does not fit the geometry

#### Finite Element Method (FEM)


Usually used in structure mechanics

- Any order of accuracy
- Based on variational methods
- Applicable on unstructured/complex grids
- Naturally implicit, well-adapted to non-linear behaviour
- Excellent for diffusion dominated problem
- Naturally implicit, need sparse linear solver (can be expensive for large systems)
- Not conservative
- More complex

#### Finite Volume Method (FVM)

Usually used in fluid mechanics

- Applicable on unstructured/complex grids
- Naturally conservative, capture discontinuities
- Easy explicit formulation (parallel computation)
- Difficult to devise stable higher order scheme
- Not very accurate for diffusion dominated problem
- More complex



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### Numerical simulation process

What does computing engineer mean?



system analysis, determination of the main phenomena, ...  
Fluid flows, wave propagation, wave breaking, impact, fluid-structure interaction, structural analysis, ...



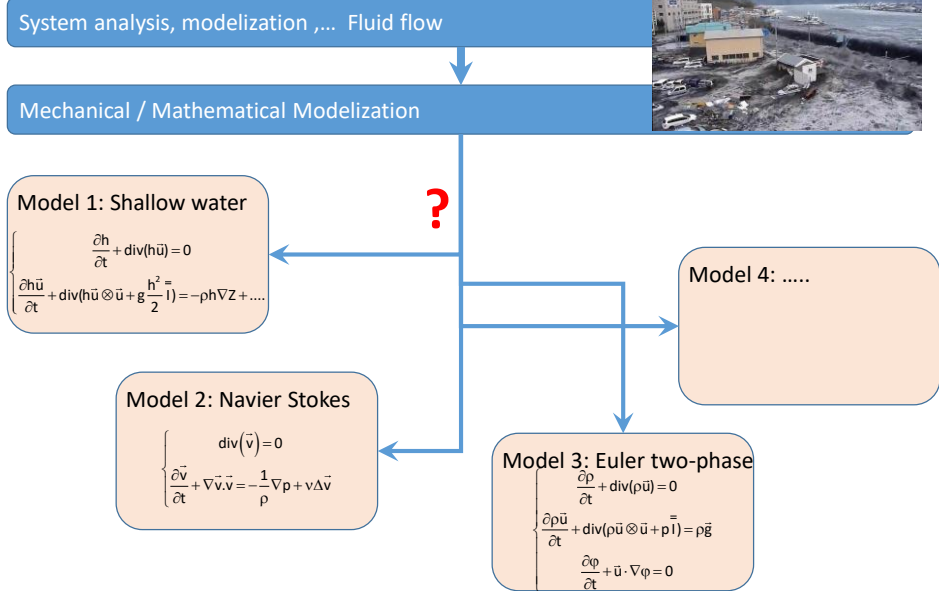
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### Numerical simulation process

What does computing engineer mean?







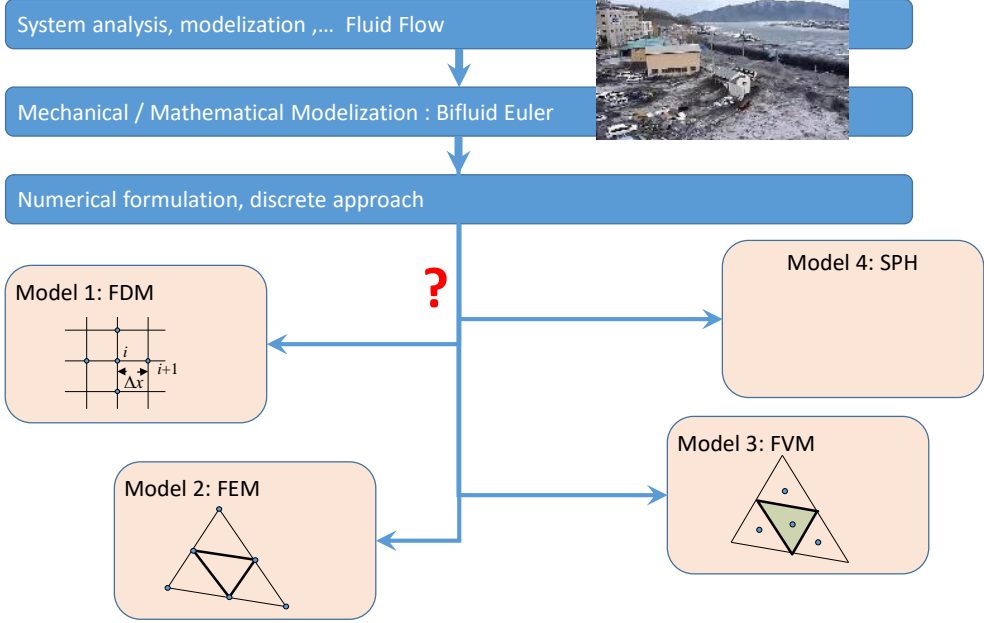

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**Numerical simulation process**

What does computing engineer mean?

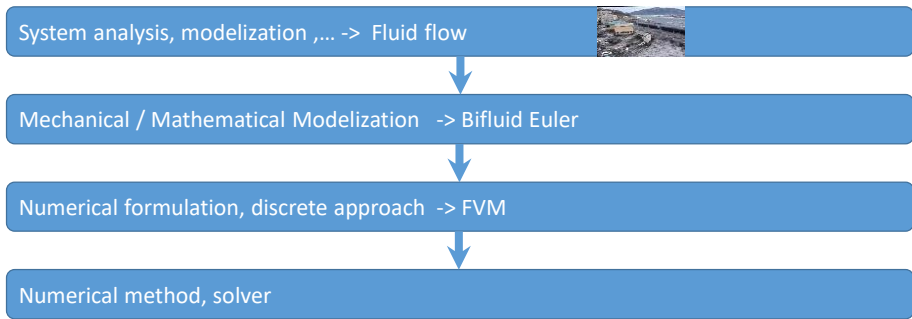
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**Numerical simulation process**

What does computing engineer mean?



- Spatial discretisation: order, gradient, ...
- Time discretization: explicit, implicit, ...
- Solver: direct, iterative, ...
- Domain decomposition, mesh refinement, ...
- Numerical improvement, ...



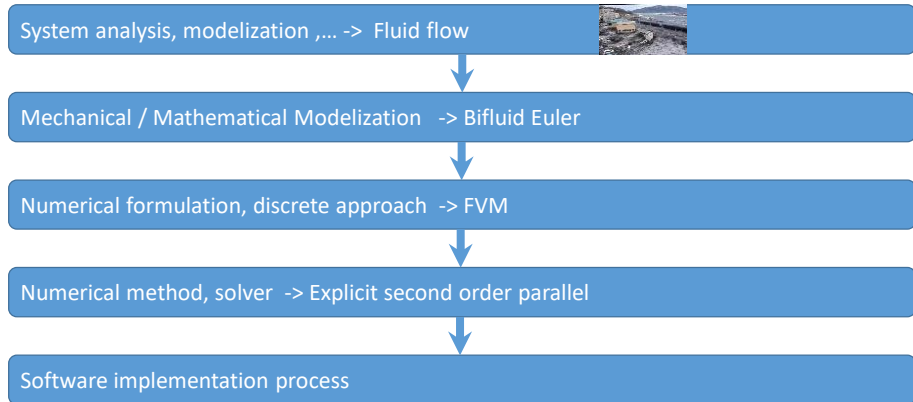
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## Numerical simulation process


What does computing engineer mean?



Programmation: C, fortran90, C++, ...

Parallel: Cuda, MPI, ...

Computer-aided software engineering, Scientific programming, versioning, ...



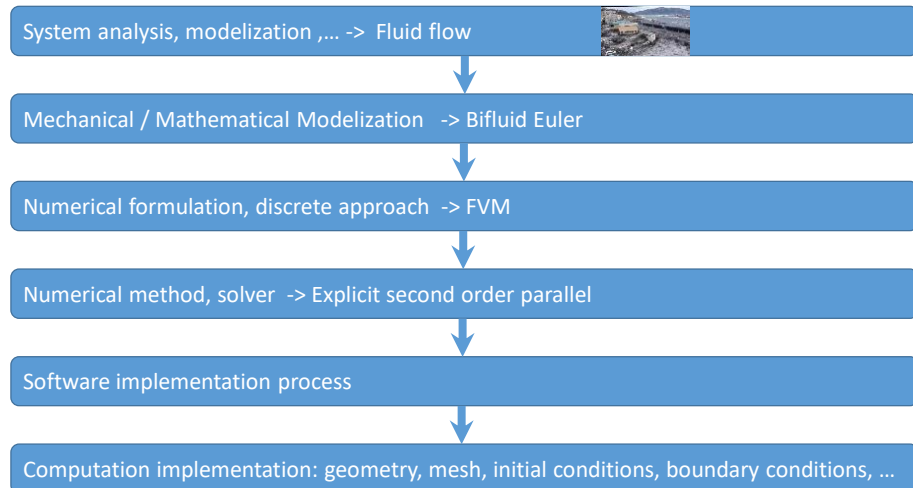
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## Numerical simulation process

What does computing engineer mean?



**Preprocessing: can be very time-consuming**

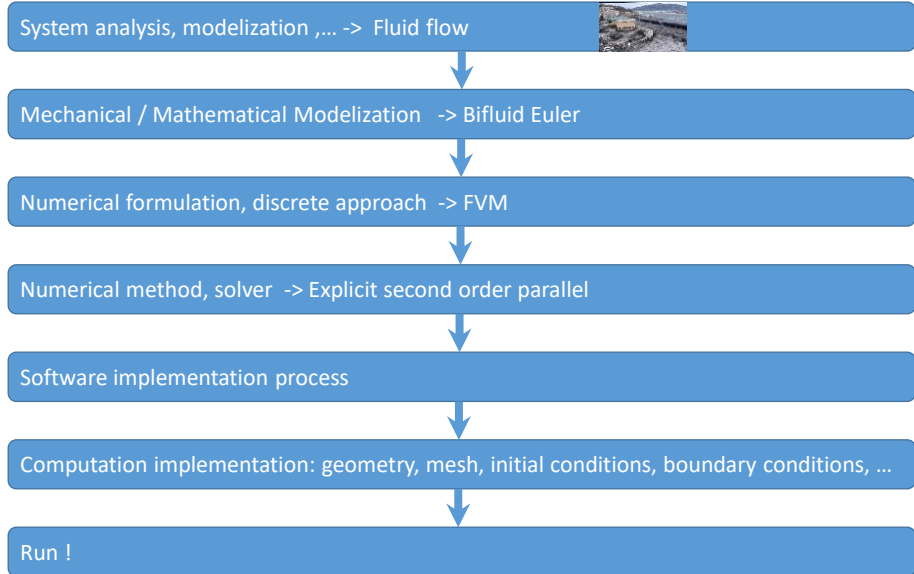



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## Numerical simulation process

What does computing engineer mean?

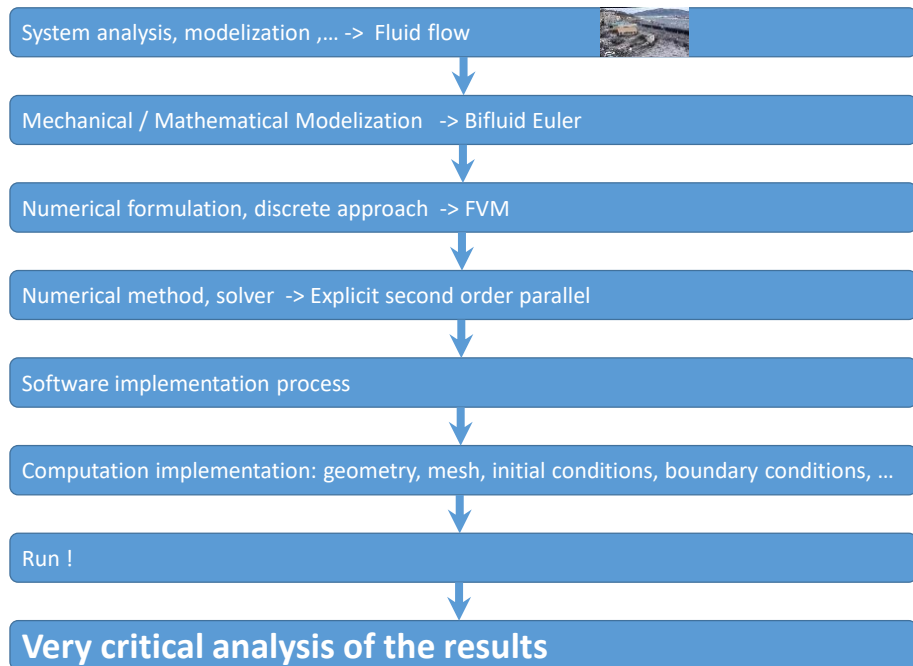



**HPC: High Performance Computing**

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## Numerical simulation process


What does computing engineer mean?

Frame 001 | 7 Jun 2004 | ts-0.000000000000000000E+000



CFD = Computational Fluid Dynamics

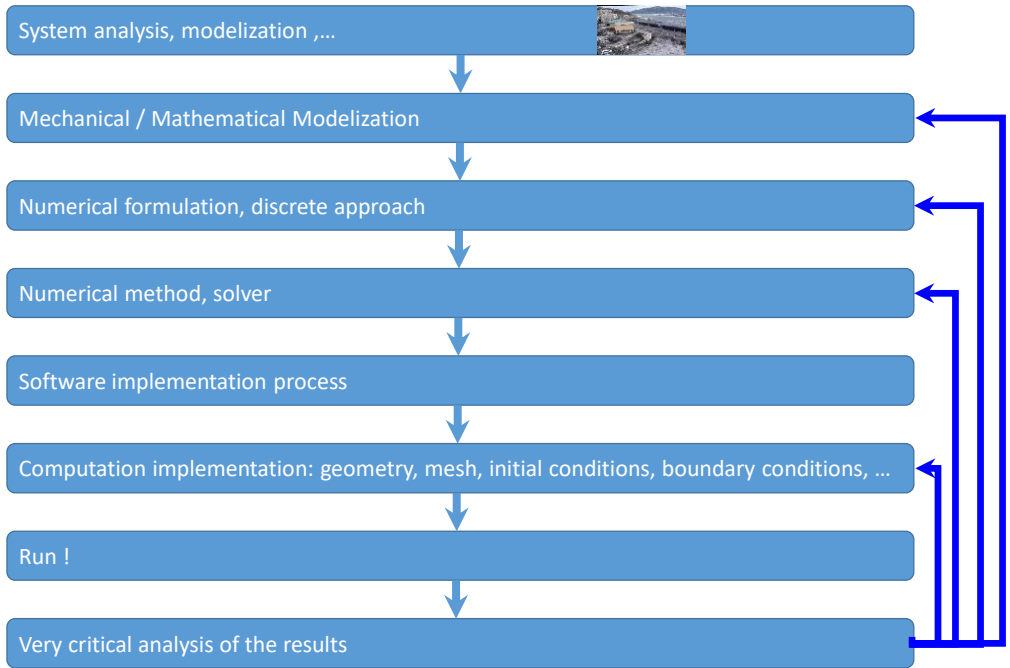
CFD ≠ Color Fluid Dynamics (G. Allaire ?)





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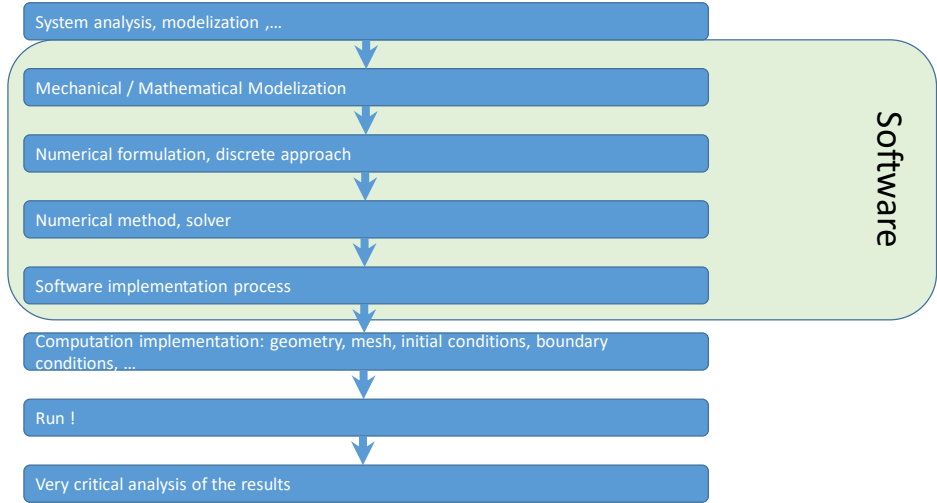

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### Numerical simulation process

What does computing engineer mean?

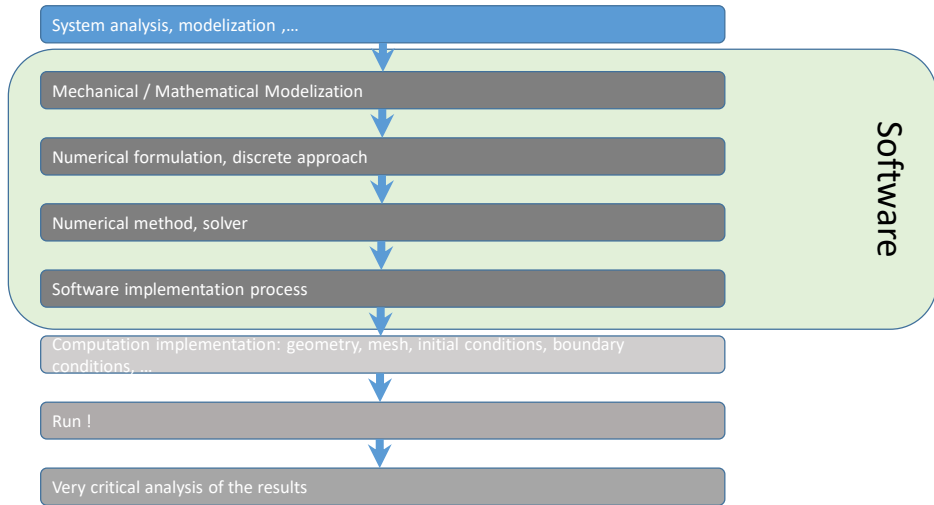
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### Numerical simulation process

What does computing engineer mean?



**Bad engineer**

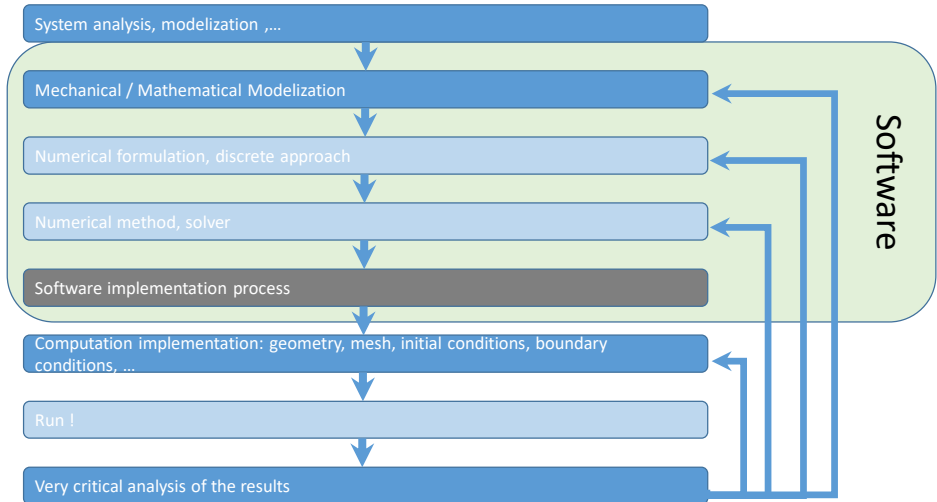


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### Numerical simulation process

What does computing engineer mean?



Most engineer

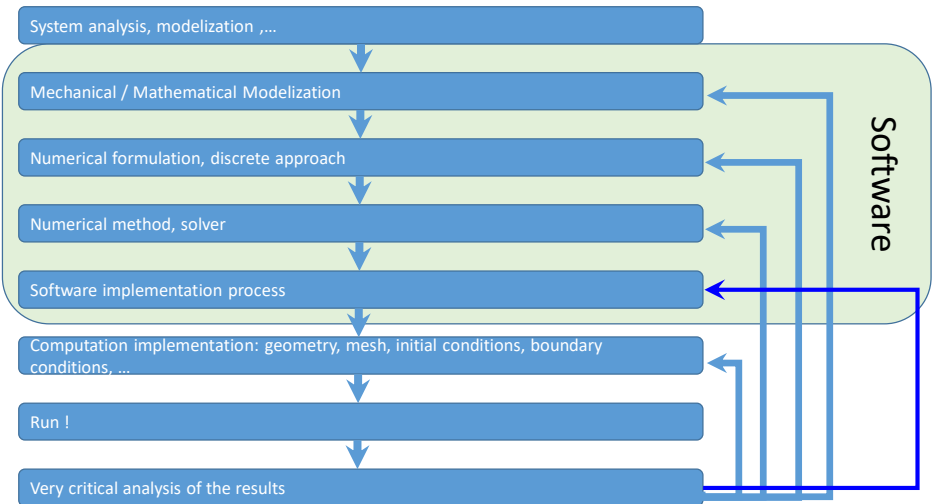


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### Numerical simulation process

What does computing engineer mean?



Engineer in computational Mechanics



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## Finite Difference Method

### Taylor's Theorem

Let  $u$  be an  $n$ -times differentiable function on an open interval containing the points  $x+h$  and  $x$ , then :

$$u(x+h) = u(x) + hu'(x) + \frac{h^2}{2!}u''(x) + \frac{h^3}{3!}u^{(3)}(x) + \dots + \frac{h^{n-1}}{(n-1)!}u^{(n-1)}(x) + \mathcal{O}(h^n)$$

Where  $\mathcal{O}(h^n) = \frac{h^n}{n!}u^{(n)}(c)$  and  $c$  is between  $x$  and  $x+h$

### Finite difference approximation of the first derivative of $u$

Forward first order finite difference approximation of  $u'$

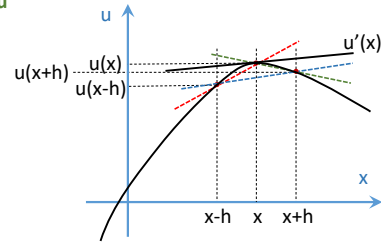

$$u(x+h) = u(x) + hu'(x) + \mathcal{O}(h^2) \quad u'(x) = \frac{u(x+h) - u(x)}{h} + \mathcal{O}(h)$$

Backward first order finite difference approximation of  $u'$

$$u(x-h) = u(x) - hu'(x) + \mathcal{O}(h^2) \quad u'(x) = \frac{u(x) - u(x-h)}{h} + \mathcal{O}(h)$$

Centered first order finite difference approximation of  $u'$

$$u'(x) = \frac{u(x+h) - u(x-h)}{2h} + \mathcal{O}(h^2)$$

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## Finite Difference Method


### Finite difference approximation of the second derivative of $u$

$$u(x+h) = u(x) + hu'(x) + \frac{h^2}{2}u''(x) + \frac{h^3}{6}u^{(3)}(x) + \mathcal{O}(h^4)$$

$$u(x-h) = u(x) - hu'(x) + \frac{h^2}{2}u''(x) - \frac{h^3}{6}u^{(3)}(x) + \mathcal{O}(h^4)$$

$$u(x+h) + u(x-h) = 2u(x) + h^2u''(x) + \mathcal{O}(h^4)$$

$$u''(x) = \frac{u(x+h) - 2u(x) + u(x-h)}{h^2} + \mathcal{O}(h^2)$$



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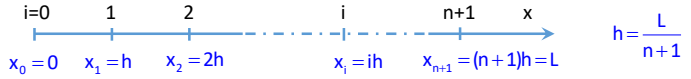
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**Finite Difference Method Elliptic equation**

**1D discretization of the Laplacian operator**


$$\Delta u = 0$$

$$\frac{\partial^2 u}{\partial x^2} = 0$$



Notation  $u(x_i) = u_i$

$$\forall i = 1, n \quad u''(x_i) \approx \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} = 0$$



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**Finite Difference Method Elliptic equation**

**2D discretization of the Laplacian operator**

$$u(x+h_x, y+h_y) = \sum_{0 \leq j+k \leq p} \frac{h_x^j h_y^k}{j!k!} \frac{\partial^{j+k} u(x,y)}{\partial x^j \partial y^k}$$

$$\Delta u(x,y) = 0$$

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

$$u(x+h_x, y) = u(x,y) + h_x \frac{\partial u(x,y)}{\partial x} + \frac{h_x^2}{2} \frac{\partial^2 u(x,y)}{\partial x^2} + \frac{h_x^3}{6} \frac{\partial^3 u(x,y)}{\partial x^3} + \mathcal{O}(h_x^4)$$


$$u(x-h_x, y) = u(x,y) - h_x \frac{\partial u(x,y)}{\partial x} + \frac{h_x^2}{2} \frac{\partial^2 u(x,y)}{\partial x^2} - \frac{h_x^3}{6} \frac{\partial^3 u(x,y)}{\partial x^3} + \mathcal{O}(h_x^4)$$

$$u(x, y+h_y) = u(x,y) + h_y \frac{\partial u(x,y)}{\partial y} + \frac{h_y^2}{2} \frac{\partial^2 u(x,y)}{\partial y^2} + \frac{h_y^3}{6} \frac{\partial^3 u(x,y)}{\partial y^3} + \mathcal{O}(h_y^4)$$

$$u(x, y-h_y) = u(x,y) - h_y \frac{\partial u(x,y)}{\partial y} + \frac{h_y^2}{2} \frac{\partial^2 u(x,y)}{\partial y^2} - \frac{h_y^3}{6} \frac{\partial^3 u(x,y)}{\partial y^3} + \mathcal{O}(h_y^4)$$

$$\frac{\partial^2 u(x,y)}{\partial x^2} + \frac{\partial^2 u(x,y)}{\partial y^2} = \frac{u(x+h_x, y) - 2u(x,y) + u(x-h_x, y)}{h_x^2} + \frac{u(x, y+h_y) - 2u(x,y) + u(x, y-h_y)}{h_y^2} + \mathcal{O}(h_x^2) + \mathcal{O}(h_y^2)$$





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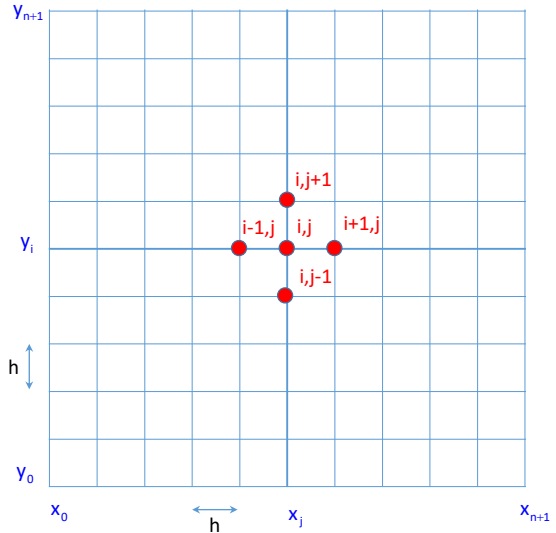
**Finite Difference Method Elliptic equation**

**2D five points stencil**

Notation  $u(x_i, y_j) = u_{i,j}$

$\forall i, j = 1, n$

$$\Delta u(x_i, y_j) \approx \frac{u_{i+1,j} + u_{i,j+1} - 4u_{i,j} + u_{i-1,j} + u_{i,j-1}}{h^2} = 0$$



**Finite Difference Method Parabolic equation**

**1D discretization**

$$\begin{cases} \frac{\partial u}{\partial t}(x, t) - \lambda \Delta u = 0 \\ x(x, 0) = u_0(x) \end{cases}$$



Notation  $u(x_i, t_n) = u_i^n$


$$\frac{\partial u}{\partial t}(x_i, t_n) \approx \frac{u_i^{n+1} - u_i^n}{\delta t} \quad \text{First order approximation}$$

**Time discretisation: explicit Euler scheme**

$$\frac{u_i^{n+1} - u_i^n}{\delta t} = \lambda \frac{\partial^2 u}{\partial x^2}(x_i, t_n) \approx \lambda \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{h^2} \rightarrow u_i^{n+1} = u_i^n + \frac{\lambda \delta t}{h^2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$$

**Time discretisation: implicit Euler scheme**


$$\frac{u_i^{n+1} - u_i^n}{\delta t} = \lambda \frac{\partial^2 u}{\partial x^2}(x_i, t_{n+1}) \approx \lambda \frac{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}{h^2} \rightarrow -u_{i+1}^{n+1} + \left( \frac{h^2}{\lambda \delta t} + 2 \right) u_i^{n+1} - u_{i-1}^{n+1} = \frac{h^2}{\lambda \delta t} u_i^n$$



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**Finite Difference Method Parabolic equation**

**Time discretisation: semi-implicit scheme**

$$\frac{u_i^{n+1} - u_i^n}{\delta t} = \theta \lambda \frac{\partial^2 u}{\partial x^2}(x_i, t_{n+1}) + (1-\theta) \frac{\partial^2 u}{\partial x^2}(x_i, t_n) \approx \theta \lambda \frac{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}{h^2} + (1-\theta) \lambda \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{h^2}$$

$$A = \begin{bmatrix} 2 & -1 & 0 & \dots & \dots & \dots & 0 \\ -1 & 2 & -1 & 0 & & & 0 \\ 0 & \ddots & \ddots & \ddots & \ddots & & \vdots \\ \vdots & \ddots & -1 & 2 & -1 & \ddots & \vdots \\ \vdots & & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & & & 0 & -1 & 2 & -1 \\ 0 & \dots & \dots & \dots & 0 & -1 & 2 \end{bmatrix} \quad U^n = \begin{bmatrix} u_1^n \\ \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \\ u_N^n \end{bmatrix} \quad \left( I_d + \frac{\theta \lambda \delta t}{h^2} A \right) U^{n+1} = \left( I_d - \frac{(1-\theta) \lambda \delta t}{h^2} A \right) U^n$$

- θ=0 explicit scheme: Forward Euler Conditionally stable
- θ=1 implicit scheme: Backward Euler Unconditionally stable
- θ=0,5 semi implicit scheme: Crank-Nicholson Unconditionally stable for θ≥0,5



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**Variational approach**

**Principle**

$$\frac{d\bar{w}}{dt} + \text{div} \bar{F}(w) = \bar{S} \quad \text{in } \Omega$$

$$\int_{\Omega} \bar{\varphi} \left( \frac{d\bar{w}}{dt} + \text{div} \bar{F} \right) dv = \int_{\Omega} \bar{\varphi} \bar{S} dv - \nabla \bar{\varphi} \dots$$

The computational domain Ω is split in several cells or elements

$$\int_{\Omega} \dots dv = \sum_e \int_{\Omega_e} \dots dv$$

$$\sum_e \int_{\Omega_e} \bar{\varphi} \frac{d\bar{w}}{dt} dv + \sum_e \int_{\Omega_e} \bar{\varphi} \text{div} \bar{F} dv = \sum_e \int_{\Omega_e} \bar{\varphi} \bar{S} dv$$

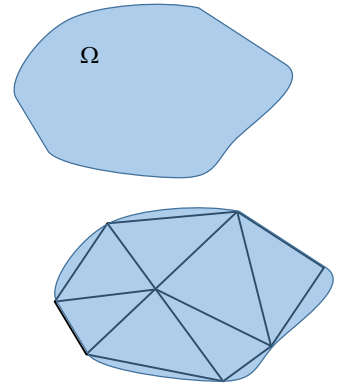
$$\sum_e \int_{\Omega_e} \bar{\varphi} \frac{d\bar{w}}{dt} dv - \sum_e \int_{\Omega_e} \nabla \bar{\varphi} : \bar{F} dv + \sum_e \int_{\Omega_e} \text{div}(\bar{F} \bar{\varphi}) dv = \sum_e \int_{\Omega_e} \bar{\varphi} \bar{S} dv$$


$$\text{div}(\bar{F} \bar{\varphi}) = \nabla \bar{\varphi} : \bar{F} + \bar{\varphi} \text{div} \bar{F}$$

$$\sum_e \int_{\Omega_e} \bar{\varphi} \frac{d\bar{w}}{dt} dv - \sum_e \int_{\Omega_e} \nabla \bar{\varphi} : \bar{F} dv + \sum_e \int_{\partial \Omega_e} \bar{F} \bar{\varphi} \cdot \bar{n} ds = \sum_e \int_{\Omega_e} \bar{\varphi} \bar{S} dv$$

Green's formula

$$\sum_e \int_{\Omega_e} \bar{\varphi} \frac{d\bar{w}}{dt} dv - \sum_e \int_{\Omega_e} \nabla \bar{\varphi} : \bar{F} dv + \sum_e \int_{\partial \Omega_e} \bar{\varphi} \bar{F} \cdot \bar{n} ds = \sum_e \int_{\Omega_e} \bar{\varphi} \bar{S} dv$$





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### Finite Element Method

#### Principle

$$\frac{dw}{dt} + \text{div} \vec{F}(w) = \vec{S} \quad \text{in } \Omega$$

$$\sum_e \int_{\Omega_e} \bar{\varphi} \cdot \frac{dw}{dt} dv - \sum_e \int_{\Omega_e} \nabla \bar{\varphi} : \vec{F} dv + \sum_e \int_{\partial \Omega_e} \bar{\varphi} \cdot \vec{F}n ds = \sum_e \int_{\Omega_e} \bar{\varphi} \cdot \vec{S} dv$$

The solutions w are approximated in each Element by a polynomial approximation

$$\vec{w}_e(x, y, z) \approx \sum_{\text{nodes of } e} N_i(x, y, z) \vec{w}_i$$

Where  $N_i$  are chosen polynomials of order 1, 2, 3, ... called interpolation functions

Rmk 1:  $N_i = 1$  on node  $i$  and :  $N_i = 0$  on other nodes

Rmk 2: the approximation of  $w$  is continuous!

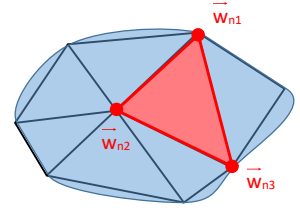
Rmk 3: the approximation of  $\nabla w$  is NOT continuous!

$$\bar{\varphi} \text{ is chosen as a continuous test function } \Rightarrow \sum_e \int_{\partial \Omega_e} \bar{\varphi} \cdot \vec{F}n ds = \sum_e \int_{\partial \Omega_e} \bar{\varphi} \cdot \vec{F}n ds \quad \text{Neumann condition}$$

Using the Galerkin's rule the test functions are chosen such that:  $\bar{\varphi}(x, y, z) \approx \sum_{\text{nodes of } e} N_j(x, y, z) \bar{\varphi}_j$

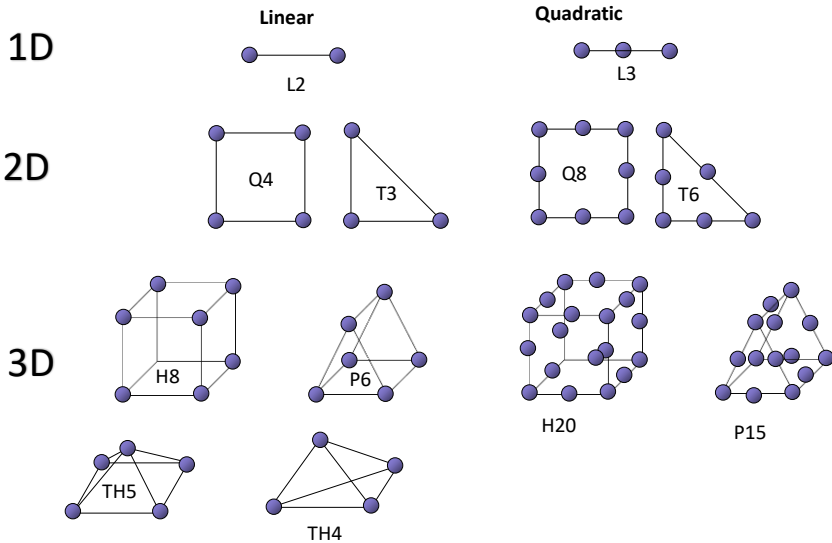

$$\sum_e \int_{\Omega_e} N_i \bar{\varphi}_j \cdot \frac{d}{dt} N_j \vec{w}_i dv - \sum_e \int_{\Omega_e} \nabla N_j \bar{\varphi}_j : \vec{F} (N_i \vec{w}_i) dv = \sum_e \int_{\Omega_e} N_j \bar{\varphi}_j \cdot \vec{S} dv - \sum_e \int_{\partial \Omega_e} N_j \bar{\varphi}_j \cdot \vec{F}n ds \quad \forall \bar{\varphi}_j$$

therefore taking successively each component of the test functions equal to 1 and the others to zero, we obtain a system of equations of size : number of degree of freedom x number of nodes



### Finite Element Method


#### Some usual finite element types

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## Finite Element Method

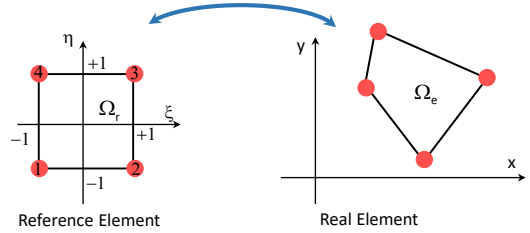
### Reference Element

$$\bar{x}(x, y, z) \approx \sum_{\text{nodes of } e} \bar{N}_i(x, y, z) \bar{x}_i \quad \bar{N}_i \text{ denote the shape functions}$$

In the case of isoparametric element, we get  $\bar{N}_i = N_i$

$$\int_{\Omega} \dots dx dy = \int_{\Omega_r} \dots \det J d\xi d\eta$$

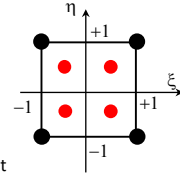

Where J denotes the Jacobian of the transformation



### Numerical integration

$$\int_{\Omega} f(\xi, \eta) d\xi d\eta = \sum_{\text{Gauss points}} \omega_{PG} f(\xi_{PG}, \eta_{PG})$$

According to the Gaussian quadrature rule,  $\omega_{PG}$  denotes the weights of each Gaussian point

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## Finite Element Method


### Many other notions specific to FEM to handle ...

- Mesh
- Conforming mesh
- Assembly
- Lumping
- Post processing analysis
- ...



### Many other notions to handle in the framework of FEM

- Boundaries conditions
- Time discretisation
- Linear solver: direct, iterative
- Non-linearities, tangent operator
- Convergence
- A posteriori error
- Domain decomposition
- Adaptive mesh refinement
- HPC
- ...



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### Finite Volume Method

#### Principle

$$\frac{d\bar{w}}{dt} + \text{div}\vec{F}(\bar{w}) = \bar{S} \text{ in } \Omega$$

$$\sum_e \int_{\Omega_e} \bar{\varphi} \frac{d\bar{w}}{dt} dv - \sum_e \int_{\Omega_e} \nabla \bar{\varphi} : \vec{F} dv + \sum_e \int_{\partial\Omega_e} \bar{\varphi} \vec{F} \cdot \vec{n} ds = \sum_e \int_{\Omega_e} \bar{\varphi} \bar{S} dv$$

The solutions  $w$  are approximated in each volume by an average value

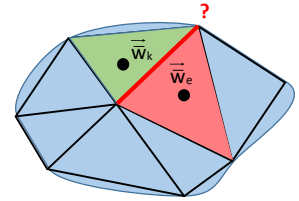
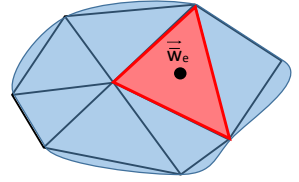

$$\bar{w}_e(x, y, z) \approx \frac{1}{|\Omega_e|} \int_{\Omega_e} \bar{w} dv$$

$\bar{\varphi}$  is chosen successively as **constant** on volume  $\Omega_e$  and null elsewhere

$$\Rightarrow \nabla \bar{\varphi} = \vec{0}$$

$$\Rightarrow \int_{\Omega_e} \frac{d\bar{w}}{dt} dv + \int_{\partial\Omega_e} \vec{F} \cdot \vec{n} ds = \int_{\Omega_e} \bar{S} dv$$

$$\Rightarrow \frac{d\bar{w}_e}{dt} + \int_{\partial\Omega_e} \vec{F}(\bar{w}_e) \vec{n} ds = \bar{S} \quad \vec{F}(\bar{w}_e) \text{ is NOT defined on } \partial\Omega_e!$$

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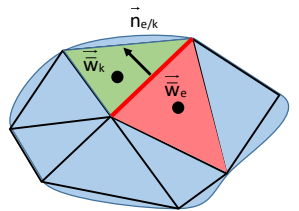
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### Finite Volume Method

#### Fluxes computation

$$\int_{\partial\Omega_e} \vec{F} \cdot \vec{n} ds \approx \sum_{\text{edges}} |\text{edge}| f(\bar{w}_e, \bar{w}_k, \vec{n}_{e/k})$$

- Centered
- Roe
- Lax-Friedrichs
- Engquist-Osher
- Lax-Wendroff
- **Godunov : exact Riemann solver**
- ...





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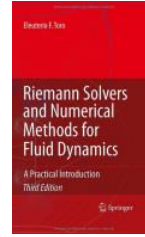

### Finite Volume Method

#### Many other notions specific to FVM to handle ...

- Structured Mesh, unstructured mesh
- Second order reconstruction (MUSCL)
- Slope limiters
- Post processing analysis
- ...

#### Many other notions to handle in the framework of FVM

- Boundaries conditions
- Time discretisation
- Linear solver: direct, iterative, multigrid
- Interface sharpening
- Numerical diffusion
- Convergence
- A posteriori error
- Domain decomposition, parallel computing
- Adaptive mesh refinement
- HPC
- ...

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### Galerkin Discontinuous Finite Element Method

#### Principle

$$\frac{dw}{dt} + \text{div}(\vec{w}) = \vec{S} \text{ in } \Omega$$

$$\sum_e \int_{\Omega_e} \vec{\varphi}_e \cdot \frac{d\vec{w}}{dt} dv - \sum_e \int_{\Omega_e} \nabla \vec{\varphi}_e : \vec{F}^T dv + \sum_e \int_{\partial\Omega_e} \vec{\varphi}_e \cdot \vec{F}n ds = \sum_e \int_{\Omega_e} \vec{\varphi}_e \cdot \vec{S} dv$$

The solutions w are approximated in each Element by a polynomial approximation of order p

$$\vec{w}_e(x, y, z) \approx \sum_{i,j,k=0}^{p-i, p-j, p-k} a_{ijk} x^i y^j z^k$$

Rmk 1: the approximation of w NOT is continuous!

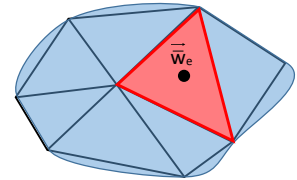
Rmk 2: We present a monomial approach, easier to implement, but harder to physically understand. A nodal approximation is still possible


$\vec{\varphi}$  is chosen successively as a **DIS**continuous test function on element  $\Omega_e$  and null elsewhere

Using the **Galerkin's** rule the test functions are also chosen such that:  $\vec{\varphi}_e(x, y, z) \approx \sum_{i,j,k=0}^{p-i, p-j, p-k} \varphi_{ijk} x^i y^j z^k$

$$\int_{\Omega_e} \vec{\varphi}_e \cdot \frac{d\vec{w}_e}{dt} dv - \int_{\Omega_e} \nabla \vec{\varphi}_e : \vec{F}^T dv + \int_{\partial\Omega_e} \vec{\varphi}_e \cdot \vec{F}n ds = \int_{\Omega_e} \vec{\varphi}_e \cdot \vec{S} dv \quad \forall \vec{\varphi}_{ijk}$$

therefore taking successively each component of the test functions equal to 1 and the others to zero, we obtain a system of equations of size : number of degree of freedom x number of coefficient of the polynomial approximation





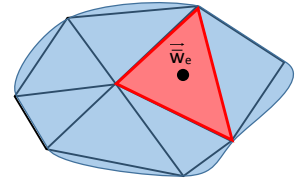
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### Galerkin Discontinuous Finite Element Method

#### Fluxes Computation

$$\int_{\Omega} \bar{\varphi}_e \frac{d\bar{w}_e}{dt} dv - \int_{\Omega} \nabla \bar{\varphi}_e : \bar{F} dv + \int_{\partial\Omega} \bar{\varphi}_e \cdot \bar{F}n ds = \int_{\Omega} \bar{\varphi}_e \cdot \bar{S} dv \quad \forall \bar{\varphi}_{ijk}$$



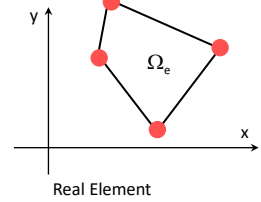
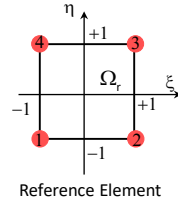
One more time, the fluxes are not defined at the interfaces between elements. An approximation has to be chosen in order to handle with the jump across the interface

#### Reference Element

$$\bar{x}(x, y, z) \approx \sum_{\text{nodes of } e} \bar{N}_i(x, y, z) \bar{x}_i$$

$$\int_{\Omega} \dots dx dy = \int_{\Omega} \dots \det J d\xi d\eta$$

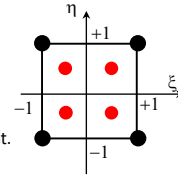

Here we use a nodal approximation, often of order 1



#### Numerical integration

$$\int_{\Omega} f(\xi, \eta) d\xi d\eta = \sum_{\text{Gauss points}} \omega_{pG} f(\xi_{pG}, \eta_{pG})$$

According to the Gaussian quadrature rule,  $\omega_{pG}$  denotes the weights of each Gaussian point. The number of Gaussian points is adapted to the order of the polynomial approximation.

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### Galerkin Discontinuous Finite Element Method

#### Many other notions specific to GD to handle ...

- Mesh, interface meshing
- Non conforming mesh
- Assembly
- Lumping
- Penalization
- Hp adaptation
- Post processing analysis
- ...



#### Many other notions to handle in the framework of GD

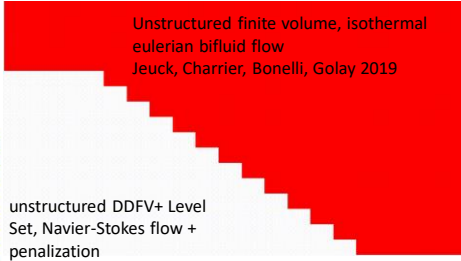
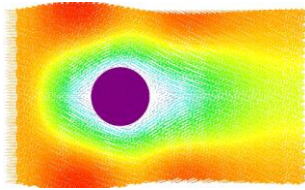
- Boundaries conditions
- Time discretisation
- Linear solver: direct, iterative
- Non-linearities, tangent operator
- Convergence
- A posteriori error
- Domain decomposition
- Adaptive mesh refinement
- HPC
- ...

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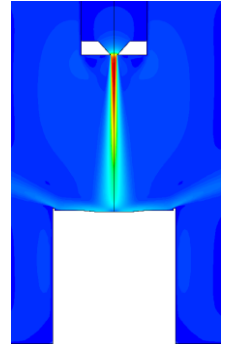
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To be continued .....



Unstructured finite volume, isothermal eulerian bifluid flow  
Jeuck, Charrier, Bonelli, Golay 2019

unstructured DDFV+ Level Set, Navier-Stokes flow + penalization  
Lakhilili, Bonelli, Golay, Galusinski, 2015



Finite volume (Fluent), Navier-Stokes flow+Turbulence+remeshing  
Mercier, Bonelli, Anselmet, 2013



structured finite volume+ Level Set, MAC scheme, Stokes flow + penalization

Lachouette, Golay, Bonelli 2009